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Experimental method for the extraction of intensity profiles by imaging the scattered pattern with a charge-coupled device

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A practicable experimental method for measuring scattering on rough surfaces is reported. The scattered patterns are captured on a screen composed of two pieces of ground glass and then imaged using a charge-coupled device. The scattered intensity profiles are extracted by converting the patterns in real space into the wave vector space. Isotropic and anisotropic samples of the rough backsides of silicon wafer are investigated respectively, and their intensity profiles are measured. The profiles of the anisotropic sample are obtained by reading the pixels on the specific orientation curves. The parameters of the samples are extracted using angle-resolved light-scattering schemes and theories. The results well agree with measurements obtained using an atomic force microscope.

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Light scattering from random rough surfaces has been a long-standing research subject. Since the 1990s, significant advancements in the characterization of rough fractal surfaces by light scattering have been achieved[1–4]. Fundamentally, Sinha et al. proposed a classical scattering model for rough self-affine fractal surfaces[5], and Yang et al. constructed mathematical relations of the half-width of the scattered intensity profiles with the parameters of rough self-affine fractal surfaces[6]. Eventually, more systematical theories and angle-resolved experimental schemes were developed for the extraction of the parameters of rough self-affine surfaces from the scattered intensity profiles[7]. Meanwhile, in the area of light scattering, scattering models and experimental methods were continuously developed to deal with the complexity of diverse practical problems[8,9]. Nowadays, light scattering has become one of the most important means of detecting surface characteristics and nanostructures. In the experimental aspect, delicate devices have been implemented for the measurement of scattered intensity profiles. One of the most frequently employed is the scatterometer designed by Shen et al.[10], which can be used to adjust scattering angles conveniently. Most works use a single detector to capture scattered light[7,10,11]. Thus, the movement of the detector makes the acquisition of an intensity profile time consuming. Another alternative device is the linear detector array[7], in which a scattered intensity profile is taken along a straight line on the measurement plane. However, for some experimental methods, light intensity has to be measured along a specific curve[10,12]. In addition, the limited extending range of the linear detector array can receive only a small part of the intensity profiles and thus induce errors in surface parameter extraction.

In this letter, we propose an experimental method for the extraction of intensity profiles by imaging the scattered pattern with a charge-coupled device (CCD) camera. The scattered pattern is first captured by a screen composed of two pieces of ground glass and then imaged by the CCD. The light intensity profiles are obtained by reading the pixels of the scattered pattern. The parameters of the two samples are obtained and compared with measurements obtained by an atomic force microscope (AFM).

The experimental setup is shown in Fig. 1(a). The optical path lies in a horizontal plane. A light beam from a He–Ne laser, which has a wavelength of 632.8 nm, is reflected by mirror M. Two polarizers, namely, P1 and P2, are used to adjust the intensity of the incident light beam. The polarization of the incident beam is kept vertical via P2. Pinhole H, which has a diameter of approximately 4 mm, is used to filter out stray light. Rough sample S is mounted on a rotational stage, with which the orientation of the sample plane and the incident angle are changed. Lens L1, with a focal length of 18 cm, is placed in front of the sample to form the Fraunhofer diffraction geometry. A screen is used to receive the scattered pattern on the focal plane of L1. The screen is composed of two pieces of ground glass, G1 and G2, and a piece of white paper clamped between them. Here, G1 and G2, having sufficient roughness (ground with 240# Al2O3 powder), and the paper piece are used so that the
For rough surfaces with Gaussian height probability distributions, the scattered intensity may be simplified from Eq. (1) as

\[ I(k_\parallel, k_\perp) = \langle U(k_\perp, k_\parallel)U^*(k_\perp, k_\parallel) \rangle = c^2 \int_{-\infty}^{\infty} \exp[-k_\perp^2 H(\rho)]\exp[-i k_\parallel \cdot d^2 \rho], \tag{4} \]

where \( H(\rho) = (|h(r_0 + \rho) - h(r_0)|)^2 \) is the height–height correlation function of the rough surface and \( c \) is a constant. For isotropic rough self-affine fractal surfaces, \( H(\rho) \) may be expressed by the phenomenological function \( H(\rho) = 2\rho^2 \{1 - \exp[-(\rho/\xi)^2]\} \) with \( \rho = \rho_0^{\text{rms}}. \)

Here, \( w \) is the root-mean-square (RMS) roughness, \( \xi \) is the correlation length of the surface, and \( \alpha \) is the roughness exponent related to the surface fractal dimension \( D_f \), which is expressed as \( \alpha = d - D_f \), with \( 0 \leq \alpha \leq 1 \) and \( d \) being the embedded dimension.

In the angle-resolved scattering scheme, the scattering profiles are taken from different incident angles, and the variation of the profiles and the specular reflection part with the change in \( k_\parallel \) are utilized. After simple calculations, Eq. (4) becomes

\[ I(k_\parallel, k_\perp) = I_s(k_\perp)\delta(k_\parallel) + I_{\text{diff}}(k_\parallel, k_\perp), \tag{5} \]

with the reflection and diffused parts denoted as \( I_s(k_\perp) = c^2\exp(-k_\perp^2 w^2) \) and

\[ I_{\text{diff}}(k_\parallel, k_\perp) = c^2 \int_{-\infty}^{\infty} \left[ \exp[-k_\perp^2 H(\rho)/2] - \exp[-k_\perp^2 w^2] \right] \cdot \exp(-ik_\parallel \cdot \rho)d^2 \rho. \tag{6} \]

Here, the expression for \( I_s(k_\perp) \) can be deduced by averaging the scattered light field \( U(k_\perp, k_\parallel) \) and using the Gaussian probability density function of surface height \( h(r_0) \), as shown by Liu et al.\[13\]. From Eqs. (5) and (6), the ratio \( R_3 \) of the integrated intensity of the specular reflection to the total integrated intensity is expressed as \( R_3 = I_s/I = \exp(-k_\perp^2 w^2) \). Then, this ratio can be used to extract the roughness \( w \) of the samples. Although the analytical expression of the transform of Eq. (6) is unattainable, approximations may be made and the half-width \( W_p \) of the diffused part of the profile may be obtained in the limit cases\[14\]

\[ W_p = \begin{cases} \frac{2}{\xi} & k_\perp^2 w^2 \ll 1 \\ \frac{2}{(k_\perp^{-1/\alpha} w^{-1/\alpha})^2} & k_\perp^2 w^2 \gg 1 \end{cases}. \tag{7} \]

With this equation and the angle-resolved experimental setup, the value of the correlation length can be obtained using \( \xi = 2/W_p \) from the half-width of the profile provided that \( k_\perp^2 w^2 \ll 1 \). The roughness exponent \( \alpha \) can be extracted from the variation of the half-widths \( W_p \) of the profile versus \( k_\perp \), with \( \ln W_p \propto (1/\alpha) \ln k_\perp \), by using the other limit case, i.e., \( k_\perp^2 w^2 \gg 1 \), and by measuring
the geometrical relation, a point \((x, y)\) on the surface are parallel to the sides of the images. The measurement, the anisotropic sample was carefully ad-
justed so that sides of the small square-pit-like structures are obtained using \(O_0 - x_0 y_0 z_0\) and \(O_{xy}\). The receiving screen in Fig. 2 is perpendicular to the direction of reflection and lies in the \(O_{xy}\) plane. Therefore, from the geometrical relation, a point \((x, y)\) on the \(O_{xy}\) plane may be expressed in the \(O_0 - x_0 y_0 z_0\) coordinates using
\[
\begin{aligned}
z_0 &= z_f \cos \theta_1 - x \sin \theta_1 \\
x_0 &= z_f \sin \theta_1 + x \cos \theta_1 \\
y_0 &= y
\end{aligned}
\] (8)

As presented in Fig. 2, the relation between the angles and the coordinates can be denoted as
\[
\begin{aligned}
\sin \theta_2 &= \frac{\sqrt{x_0^2 + y_0^2}}{\sqrt{x_0^2 + y_0^2 + z_0^2}} \\
\sin \theta_3 &= \frac{y_0}{\sqrt{x_0^2 + y_0^2}} \\
\cos \theta_3 &= \frac{x_0}{\sqrt{x_0^2 + y_0^2}}
\end{aligned}
\] (9)

With Eqs. (8) and (9), as well as Eqs. (2) and (3), the scattered intensity distributions on the receiving screen can be converted into intensity profiles with variables \(k_\parallel\) and \(k_\perp\). Afterward, the surface parameters may be extracted.

Next, an experiment on the two samples of the backsides of commercial silicon wafers is conducted. The samples are measured using an AFM, and their morphological images are shown in Figs. 3(a) and (b), respectively. Four images from different areas of each sample were taken to obtain average data and hence avoid errors due to the limited scan range. For the AFM measurement, the anisotropic sample was carefully adjusted so that sides of the small square-pit-like structures on the surface are parallel to the sides of the images. The roughness values \((w)\) of the samples are calculated using the height data from their AFM images and are given in Table 1. The averaged height–height correlation function \(H(\rho)\) of the images is calculated to determine the lateral correlation length \(\xi\) and roughness exponent \(\alpha\) from the AFM images. The roughness exponent \(\alpha\) is obtained by linearly fitting the height–height correlation function curve with the \(\rho < \xi\) region, where \(\ln[H(\rho)] \propto 2 \alpha \ln(\rho)\). On the other hand, the values of \(\xi\) are obtained by fitting the numerical function \(H(\rho)\) with the phenomenological function of the height–height correlation for the self-affine fractal model.

In the light-scattering measurement, each image of the scattered patterns at a certain angle of incidence is sampled 100 times with the CCD accumulations to reduce electric noise. In addition, four images of the pattern are taken at different positions of the same sample to obtain averaged data and eliminate speckles. The disturbance of the speckle effect on the scattered profiles can be decreased by using an imaging lens with a large relative aperture \((f/1.8)\) to produce speckles with extremely small average sizes such that a CCD pixel can cover several (approximately over 10) speckle grains. Then, the effect of speckles is physically averaged by the CCD detection. The scattered patterns of each sample are taken at incidence angles from 40° to 80°, with increments of 5°, and at 83°. Some of scattered pattern images for the two samples are shown in Figs. 4(a) and (b).

The differences among the scattered patterns of the isotropic and anisotropic samples and the evolutions of the patterns with the angle of incidence can be clearly seen in Fig. 4. For the isotropic sample, the distribution of the intensity pattern is nearly round and uniform at a smaller incident angle. As the angle is increased, the pattern contracts, more evidently in the vertical direction, and appears to be oblate. For the anisotropic sample, secondary peaks exist on the four sides of the main peak. As the incident angle is increased, the lower and upper

![Fig. 3. AFM images of the two kinds of rough surface samples. (a) Isotropic sample with 50×50 (μm) scan area and (b) anisotropic sample with 100×100 (μm) scan area.](image)

<table>
<thead>
<tr>
<th>Sample</th>
<th>(w) (μm)</th>
<th>(\xi) (μm)</th>
<th>(\alpha)</th>
</tr>
</thead>
<tbody>
<tr>
<td>AFM</td>
<td>LS</td>
<td>AFM</td>
<td>LS</td>
</tr>
<tr>
<td>Isotropic</td>
<td>0.47±0.03</td>
<td>0.34</td>
<td>4.16±0.17</td>
</tr>
<tr>
<td>Anisotropic</td>
<td>0.74±0.10</td>
<td>0.63</td>
<td>12.07±0.10</td>
</tr>
</tbody>
</table>
The scattered intensity profiles of the isotropic sample are extracted from the intensity patterns by reading the intensity data in the pixels along the orientation curves, as shown in Fig. 4(a). Along each line segment, the coordinate \( y \) in \( Oxy \) varies as in Fig. 2, whereas \( \theta_2 \) changes insignificantly. Based on Eqs. (2), (3), (8), and (9), \( k_\perp \) therefore changes with \( y \), and \( k_\parallel \) can be approximately regarded as a constant. Then, with the coordinates converted into \( k_\parallel \), the intensity profile \( I(k_\perp, k_\parallel) \) at an incident angle is obtained. Figure 5(a) presents the typical normalized profile curves at incident angles of 40°, 55°, 70°, and 75°, respectively. The parameters of the sample can be obtained using the aforementioned method. In the extraction of the roughness exponent, the ratio \( R_3 \) is calculated with the profile at 83° to obtain \( w=0.63 \) \( \mu \)m. The lateral correlation length is extracted from the profile at 83° as \( \xi=11.36\pm0.46 \) \( \mu \)m. The roughness exponent is obtained as \( \alpha=0.77 \pm0.09 \) from the half-widths of the profiles at 45°, 50°, 55°, and 60°. Table 1 shows the comparison between these parameters and the AFM measurements.

Table 1 shows that the results obtained by the light-scattering method roughly conform with the AFM measurements. In general, the roughness values obtained by light scattering are smaller than those obtained by the AFM. Such discrepancy may be due to the deviation in height as the AFM measurements have a relative tilt height background. Other possible reasons for the discrepancy in surface parameters may be factors such as the approximations taken in the theoretical analysis, the systematic errors in the experiments, and the conversion of coordinates into the wave vector space.

In the experiment, the central \( \delta \) peak should be measureable, which would limit the roughness of the samples in a suitable range. Practical experience shows that a roughness value from 0.05 to 1.20 \( \mu \)m is sufficiently accurate. The lateral correlation length \( \xi \) determines the half-width \( W_p \) of the intensity profile in the case of \( k_\perp^2 w^2 \ll 1 \). Thus, the feasibility of resolving the whole scattered profile determines the measurable range. According to experience, this range is approximately from 2.0 to 15.0
µm. In addition, several assumptions and approximations were used in the theoretical derivations. Firstly, the Kirchhoff theory of diffraction under the condition of Eq. (1) is widely used in literature and is shown to be sufficiently accurate. Secondly, for the surface model, a self-affine fractal rough surface with Gaussian distribution was used, and based on which, Eq. (4) and the $H(\rho)$ expression were obtained. Furthermore, the AFM data of the samples were analyzed, and the results showed good consistency with those of the model. For the extraction of roughness exponent $\alpha$, the condition $k_\perp^2 w^2 \gg 1$ in Eq. (7) was used, which may be evaluated by the linear proportion of $W_p$ versus $\ln k_\perp$ in the large $k_\perp$ region, and accurate results were obtained.

In conclusion, we present a convenient and practicable method for the light-scattering measurement of rough surface samples, which is done by imaging the scattering patterns and using light-scattering theories. Aside from providing clear images for the study of the scattering properties of the samples, the experimental setup also reduce the difficulty in the acquisition of scattering intensity data. The convenience in reading the data in any pixels of the images facilitated the analysis of the scattering of the anisotropic sample by taking the profiles along the specific orientation curves. The experimental results indicate that the setup and method proposed in this letter are effective and convenient for the light-scattering measurement of random rough surfaces.

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References