Incoherent Fourier ptychographic photography using structured light

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1. INTRODUCTION

Structured light illumination is important in computational photography and image-based rendering. It has been widely used in object recognition [1,2], 3D shape recovery [3], 4D light field rendering [4], synthetic aperture confocal imaging [5], and dual photography [6]. For example, Microsoft Kinect uses an IR projector to project a structured dense light pattern on the sample and a postprocessing algorithm to recover the 3D image of the object. A 3D laser light scanner projects a pattern of parallel stripes on the object and recovers the 3D information based on the geometrical deformation of the strips. In dual photography [6], a video projector is used to generate multiplexed patterns for sample illumination. The acquired information is then used to recover the transport matrix describing how light from each pixel of the projector arrives at each pixel of the camera. This matrix can be used to generate a new photorealistic image from the point of view of the projector.

The above-mentioned research directions and applications share a similar strategy on the system design: projecting multiple light patterns onto the sample and recovering new information based on the acquired images. However, in these approaches, the frequency mixing between the object and the illumination pattern has not been considered in the reconstruction process, and the image resolution is determined by the cutoff frequency of the employed photographic lens. In this article, we propose an incoherent photography imaging approach, termed Fourier ptychographic photography (FPF), that uses the frequency mixing effect for super-resolution imaging. The concept of the reported approach is similar to the structured illumination microscopy (SIM) approach, where a sinusoidal pattern is used to modulate the high-frequency sample information into the low-frequency passband [7]. In our implementation, we project a number of high-frequency random patterns onto the sample and acquire the corresponding images of the object. Similar to SIM, the frequency mixing between the sample and the random illumination pattern shifts the high-frequency component to the passband of the collection optics. Therefore, each raw image contains information that is beyond the cutoff frequency of the photographic lens. Based on multiple images acquired under different structured light patterns, we used the Fourier ptychographic algorithm to recover both the high-resolution object image and the unknown illumination pattern. Similar to SIM, the resolution of the recovered image using the reported approach is better than the resolution limit of the lens’s aperture.

We note that the use of the frequency mixing effect for super-resolution imaging is not a new idea. This concept is well known in the microscopy community [7]. It has also been used in telescope settings [8,9]. In this paper, we will use such a frequency mixing effect for super-resolution photographic imaging. In particular, we will demonstrate, for the first time to our knowledge, the use of the Fourier ptychographic recovery scheme for the photographic imaging setting. We will show that the reported approach is able to recover both the super-resolution image of the object and the unknown illumination pattern at the same time. Since the system design of the reported approach is compatible with many illumination-based imaging platforms, it may provide new insight...
for the development of computational photography and find applications in remote sensing, active-illumination night vision systems, and imaging radar.

In the following, we will first introduce the principle and the setup of the reported approach. We will then demonstrate the imaging performance using various targets. Finally, we will discuss the advantages and limitations of the reported approach.

2. FOURIER PTYCHOGRAPHIC PHOTOGRAPHY

Figure 1 shows the schematic of our FPP setup. In the illumination path, we projected the image of a semitransparent diffuser (white paint sprayed on a glass slide) onto the object. The speckle size of the projection pattern is in the range from 200 to 500 μm. In the detection path, we used a CCD and a photographic lens (Nikon, 50 mm) to capture the image of the object. The f-number of the lens was adjusted to 22, with the optical resolution matching the speckle size of the projection pattern. We then used a mechanical scanner to move the semitransparent diffuser to different positions, and thus the corresponding projection pattern shifts across the object (Media 1). For each projection pattern $P_n (n = 1, 2, 3, ...)$, we acquired one corresponding image of the object $I_n (n = 1, 2, 3, ...).$ Based on all acquired images, we recovered the high-resolution sample image $I_{obj}$ and the unknown projection pattern $P_n.$ Similar to SIM, the final resolution of $I_{obj}$ bypasses the limit of the collection optics.

In our implementation, we used the iterative Fourier ptychographic algorithm to recover both $I_{obj}$ and $P_n.$ The Fourier ptychography (FP) technique [10–16] is a recently developed approach that facilitates high-resolution imaging beyond the diffraction limit of the employed optics. This FP approach shares its roots with phase retrieval techniques such as the Gerchberg–Saxton algorithm [17], ptychography [18], and other error-reduction phase retrieval methods [19,20]. The recovery process of FP iteratively switches between two working domains: the spatial and Fourier domains. In the spatial domain, the acquired image is used as the object constraint for the solution. In the Fourier domain, the optical transfer function (OTF) of the objective lens is used as the support constraint for the solution. The recovery process of the FP algorithms is summarized in Fig. 2. For the case of incoherent FP [16], it starts with an initial guess of the $I_{obj}$ and $P_n.$ In each iteration step, the initial guess is updated by the acquired intensity image $I_n$, in both the spatial and Fourier domains. The iteration is repeated until the solution converges (it stops if the difference of the solutions from sequential iterations is less than a predefined value).

The updating process in the Fourier domain (step 3, Fig. 2) can be expressed as follows:

$$\mathcal{F}(I_{tn}^{\text{updated}}) = \mathcal{F}(I_{tn}) + \text{OTF} \cdot (\mathcal{F}(I_n) - \text{OTF} \cdot \mathcal{F}(I_{tn})).$$

where $\mathcal{F}()$ denotes Fourier transform and OTF denotes the optical transfer function of the employed lens. Unlike with the conventional deconvolution approach, Eq. (1) does not amplify the noise of the image, because the OTF is not a denominator. The updating process in the spatial domain can be expressed as follows (step 4):

$$I_{obj}^{\text{updated}} = I_{obj} + \frac{P_n}{(\text{max}(P_n))^2} \cdot (I_{tn}^{\text{updated}} - I_{obj} \cdot P_n).$$

The super-resolution information comes from the multiplication between the illumination pattern and the object in Eq. (2).

In our implementation, we translated one unknown illumination pattern to different spatial positions, $x_n$, and captured the corresponding images of the object. Following the same logic of Eq. (2), we can recover the unknown illumination pattern $P_{\text{unknown}}$ as follows (the initial guess of $P_{\text{unknown}}$ can be a constant number):

$$P_{\text{unknown}}^{\text{updated}} = P_{\text{unknown}} + \frac{I_{obj}}{(\text{max}(I_{obj}))^2} \cdot (I_{tn}^{\text{updated}} - I_{obj} \cdot P_{\text{unknown}}(x - x_n)).$$

where $x_n$ ($n = 1, 2, 3, ...$) represents the translational positions of the illumination pattern. These positions were controlled by the motion stage in our platform, and they were known information. If the shifted positions of the illumination pattern are unknown, it is possible to use the cross correlation of the acquired images to recover them [21].

The iterative updating process using Eqs. (1)–(3) is inspired by the extended ptychographic iterative engine developed by Maiden and Rodenburg [22]. It is the same as the pupil recovery scheme of the coherent FP approach [15]. The only
difference is the incoherent intensity patterns versus angle-
varied coherent planes waves.

3. IMAGING PERFORMANCE OF THE FPP

To demonstrate the imaging performance of the reported FPP approach, we first used a resolution target as the sample and followed the FPP procedures to recover the super-resolution image.

Figure 3 shows the result using the reported FPP platform. Figure 3(a) shows the reference image captured under uniform illumination, (b1) the captured raw image under pattern illumination, (b2) the recovered image using 100 raw images, (b3) the recovered illumination pattern, (b4) line traces of (a) and (b2). Also refer to Media 1.

We also used the reported platform to image different objects, and the results are shown in Fig. 5. Figures 5(a1)–5(c1) show the captured images under uniform illumination (also refer to Media 2). Figures 5(a2)–5(c2) show the recovered images using 100 raw images and 15 loops for the reconstruction process. In particular, we note that we are not solving the pixel aliasing problem using a subpixel shift in the raw images. The super-resolution effect here is to bypass the diffraction limit of the photographic lens. The raw image we captured is oversampled. We have shown the Fourier spectrum of

![Image](image_url)

Fig. 3. Imaging performance of the reported FPP platform. (a) The reference image captured under uniform illumination, (b1) the captured raw image under pattern illumination, (b2) the recovered image using 100 raw images, (b3) the recovered illumination pattern, (b4) line traces of (a) and (b2). Also refer to Media 1.

![Image](image_url)

Fig. 4. Image reconstruction using different numbers of raw images. The solution converges with the 16 raw images. We used 15–20 loops in this experiment.

![Image](image_url)

Fig. 5. Demonstration of the reported platform for different objects: a dollar bill, a quick response code, and an insect. (a1)–(c1) The reference images under uniform illumination, (a2)–(c2) the recovered images using the reported platform. We used 100 raw images and 15 loops for the reconstruction. Also refer to Media 2.
the raw image in the inset of Fig. 5(c1). The yellow circle denotes the cutoff frequency of the OTF of the photographic lens.

In Fig. 6, we also demonstrate the reported platform for color imaging. We used R/G/B LEDs to capture raw images at different channels and combined the reconstructions to form a high-resolution color image. The resolution improvement is evident when we compare Fig. 6(d) to Fig. 6(b).

4. CONCLUSION AND DISCUSSION

In conclusion, we have demonstrated a photographic imaging approach that uses computational illumination for super-resolution imaging. Our setup is similar to many existing computational illumination platforms, where structured light patterns are projected onto the sample and corresponding images are used to recover the hidden information of the object. In the reported approach, we used the structured light pattern to modulate the high-frequency component of the object into the low-frequency passband. We then used the Fourier ptychographic algorithm to recover the object image. We show that the reported approach is able to improve the resolution beyond the resolution limit of the photographic lens.

The design focus of conventional photographic imaging platforms is the collection optics. More and more lens elements are used in the optical design to correct for aberrations. This effort is required mainly due to the use of a 2D planar image sensor. If the sensor surface itself is curved instead of the lens elements, image capture solutions with large apertures can easily be realized by means of a simple lens system with a lower number of optical elements. However, an image sensor with a curved surface is not compatible with the existing workflow in the semiconductor industry. In the reported platform, on the other hand, we can make a semitransparent diffuser on a curved surface. Therefore, we can project high-frequency speckle patterns onto the object with a simple large-aperture lens system. We can then capture the corresponding images and use the iterative algorithm to recover the high-resolution image of the object.

There are also three limitations associated with the reported platform. (1) It is computationally intensive. Real-time reconstruction needs to be implemented using a graphical processing unit. (2) Multiple frames are needed for high-resolution reconstruction. One future direction is to design an optimal projection pattern to minimize the number of acquisitions. (3) We assume the object is static in our implementation. If the object moves during acquisition, the motion of the object may provide a mean of mechanical scanning (for example, products on a conveyer belt). In this case, we may need to track the motion of the object instead of projecting the translated patterns.

Finally, we reiterate that the use of the frequency mixing effect for super-resolution imaging is well known in the microscopy community. In this paper, we have demonstrated, for the first time to our knowledge, the use of FP for incoherent photographic imaging settings. The results of this paper may provide some insights for the computational photography and computer vision communities. It may also find applications in remote sensing, active-illumination night vision systems, and imaging radar.

REFERENCES