Observation of intracavity electromagnetically induced transparency in Cs vapor coupled with a standing-wave cavity

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The cavity transmission spectrum is experimentally investigated in A-type three-level atoms coupled to a standing-wave cavity system. It is shown that the dark-state polariton peak is not generated at resonance but rather at off-resonance. The theoretical analysis reveals that the absence of an on-resonance dark-state polariton peak is mainly caused by the strong absorption of the intracavity medium to the probe cavity mode counterpropagating with the coupling field due to the Doppler shift in the hot atoms. Moreover, the optimal frequency position of the cavity mode for an efficient dark-state polariton peak is also demonstrated.

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Cavity quantum electrodynamics (QED) is mainly researching the interaction process with a coherent atomic medium placed inside an optical resonant cavity, and has been of great interest in recent years. A well-known cavity-QED effect is the vacuum Rabi splitting or normal-mode splitting phenomenon that is under the strong coupling condition, where the cavity transmission spectrum exhibits a double-splitting-peaked profile when two-level atoms are coupled into a single cavity mode. Based on the above, when strong coherent light is injected into the intracavity medium to form a A-type three-level structure in which the electromagnetically induced transparency (EIT) is prepared, the cavity transmission spectrum produces an additional narrow central peak under the condition of the two-photon resonance, that is the intracavity EIT (i.e., dark-state polariton). Possible applications of intracavity EIT include long-lived storage of quantum information, preparation of squeezed or entangled states, and high-resolution spectroscopy. The three-peak spectrum of a three-level EIT atom-cavity system has been demonstrated in cold atoms, single atoms, hot atomic vapor, and ion Coulomb crystals. The EIT effect was also achieved with a terahertz frequency in a waveguide cavity and meta-atoms. Recently, the four-wave mixing effect based on the atom-cavity system has also been researched. However, the spectrum property of the hot three-level atom standing-wave cavity (SC) system has not been reported yet.

As the resonant probe mode is copropagating with the coupling field for the composite system with hot A-type three-level atoms inside a ring cavity (RC), which is Doppler free for the intracavity moving atoms to the probe mode and coupling frequency, the intracavity EIT is achieved at the center of atomic resonance when the two-photon resonance is satisfied. In this Letter, we experimentally demonstrate the cavity transmission spectrum in a hot Cs vapor coupled to a near-confocal standing-wave cavity (SC) by controlling the frequency detunings of the coupling field and the cavity-mode field. Quite different from the situation of an optical RC, however, it is shown that the dark-state polariton peak is created not at the center of atomic resonance but rather at off-resonance. Theoretically, the Doppler-shift effect due to the hot atoms and the coherent pump effect of the coupling fields can give a qualitative interpretation. Furthermore, the multipoint profiles and their position distributions relevant to the frequency detunings of the cavity mode and the coupling field are also experimental investigated.

The energy levels and experimental setup are shown in Fig. 1. The hyperfine structure in the D1 line of the $^{133}$Cs atom is used for the A-type three-level EIT system composed of two lower states $|a\rangle (6^2S_{1/2}, F = 4)$, $|b\rangle (6^2S_{1/2}, F = 3)$, and a common upper state $|c\rangle (6^2P_{1/2}, F = 4)$ [see Fig. 1(a)]. The coupling and probe fields drive the transitions $|a\rangle \leftrightarrow |c\rangle$ and $|b\rangle \leftrightarrow |c\rangle$ with frequencies $\omega_p$ and $\omega_c$ and detunings $\Delta_p = \omega_p - \omega_{ca}$ and $\Delta_c = \omega_c - \omega_{cb}$, respectively. $\omega_{ca}$ and $\omega_{cb}$ are the corresponding resonance frequency. A near-confocal optical SC is composed of an input mirror M1 and an output mirror M2 that have the same transmissivity ($T_{M1} = T_{M2} = 0.5\%$) and the curvature radius ($\rho_{M1} = \rho_{M2} = 150\ mm$). A piezoelectric transducer (PZT) mounted on M2 is used for scanning and locking the cavity length. The Cs vapor cell is 75 mm long with antireflection (AR) coated end windows and is wrapped in $\mu$-metal sheets for magnetic field shielding and heat tape for controlling its temperature. The vertical-polarized coupling field is injected through a polarization beam splitter
2 plots the experimentally measured cavity transmission spectrum versus the probe frequency detuning $\Delta_p$. When there is no coupling work ($P_c = 0$), the cavity transmission shows a three-peak profile, which includes two normal splitting modes and the residual resonant transmission shows a three-peak profile, which includes the cavity-mode detuning $\delta_q$, which is defined as the $q$th cavity-mode detuning [see the red curve (I) in Fig. 2]. When the coupling beam is injected and its power is $P_c = 10$ mW, the splitting of the two near-symmetric normal side peaks becomes larger because the interaction strength of atom-cavity increases from $\sqrt{g^2N}$ to $\sqrt{g^2N + \Omega_c^2/4}$, where $g$ is the single-atom-cavity coupling strength, $N$ is the number of atoms in the cavity, $\Omega_c$ is the Rabi frequency of the coupling beam. Quite different from the results of the RC[37], however, the dark-state polariton peak is not seen in the resonant case ($\delta_q = \Delta_c = 0$) [see the blue curve (II) in Fig. 2]. This is mainly because the cavity mode is recirculating in the SC, rather a strong absorption not transparency for the probe cavity mode reflected by M2 counterpropagating with the coupling beam, restrains the creation of the dark-state polariton peak [see Fig. 1(b)]. The relative discussion is seen at the next section. As the cavity transmission is detected by scanning the probe frequency with a fixed cavity length, we use the saturated absorption spectrum (SAS) to judge the frequency scanning range of the probe [see the gray curve (III) in Fig. 2]. Furthermore, the little peak of the light gray curve (IV) is the EIT signal obtained by detecting the transmission of the probe acting with the Cs vapor coupled by the co-propagating coupling field in free space (not plotted in the experimental setup in Fig. 1), which is generated under the condition of two-photon resonance ($\Delta_c = \Delta_p$) and can be used to monitor the frequency detuning of the coupling beam.

Theoretically, much attention is focused on why the dark-state polariton peak can be seen at a certain coupling detuning not close to the center of atomic resonance for a composite SC. So, the susceptibility of the intracavity atoms should be considered primarily under the free space condition. See Fig. 1(a) for a $\Lambda$-type three-level system, the density-matrix equations of motion in the rotating frame are given by[35]

$$
\rho_{ab} + \rho_{bb} + \rho_{cc} = 1,
\rho_{aa} = \gamma_c \rho_{cc} + b(\rho_{bb} - \rho_{aa}) - i\Omega_p (\rho_{ac} - \rho_{ca}),
\rho_{bh} = \gamma_c \rho_{cc} + b(\rho_{aa} - \rho_{bb}) - i\Omega_c (\rho_{bc} - \rho_{cb}),
\rho_{ca} = -(i\Delta_p + \Gamma_c) \rho_{ca} + i\Omega_c (\rho_{bc} - \rho_{cb}),
\rho_{cb} = -(i\Delta_c + \Gamma_c) \rho_{cb} + i\Omega_c (\rho_{bc} - \rho_{cb}),
\rho_{ac} = (i\Delta_c - i\Delta_p - \Gamma_c + \Gamma_ab) \rho_{ac} + i\Omega_p (\rho_{cc} + \rho_{bc}) - i\Omega_c (\rho_{ac} + \rho_{ca}),
$$

where $\gamma_c$ and $\gamma_ab$ are the decay rates from level $|c\rangle$ to $|a\rangle$ and $|b\rangle$, respectively; $\Omega_c$, and $\Omega_p$ are the Rabi frequencies of the probe and coupling fields, respectively; the off-diagonal decay rates are $\Gamma_{ij} = \gamma_{ij} + \gamma_{ji}$, which is the dephasing rate of level $i$; and $b$ is the population relaxation rates between $|a\rangle$ and $|b\rangle$.

It is expressed by solving the Eqs. (1) for $\rho_{ca}$ in a steady state

$$
\rho_{ca} = \frac{i\Omega_p((-i\Delta_c + \Gamma_{ca})(i(\Delta_c - \Delta_p) + \Gamma_{ab})\rho_{bb} - \rho_{cc}) + \Omega_c^2)}{(-i\Delta_c + \Gamma_{ca})(i(\Delta_c - \Delta_p) + \Gamma_{ab})(i(\Delta_p + \Gamma_{ca}) + \Omega_c^2)},
$$

considering the Doppler-broadening effect of the intercavity atoms, the susceptibility of the system for probe laser $\chi$ expresses the $\rho_{ca}$ as a function of the Doppler shift $v$, and integrate over the Maxwell-velocity distribution $f(kv)$,

$$
\chi = \int_{-\infty}^{\infty} d(kv)k \frac{\rho_{ca}}{\Omega_p} f(kv),
$$
where $f(kv) = (ku/\sqrt{2})^{-1} \exp[-(kv)^2/(ku)^2]$, $k$ is the wave vector, $u = \sqrt{2k_BT/m}$ is the most probable speed of an atom at a given temperature $T$ and atomic mass $m$, $k_B$ is the Boltzmann constant, $h = N\sigma^2/ε_0\hbar$ is a constant, and $N$ is the atomic density at temperature $T$.

In Fig. 1(b), as the probe beam coupled to the standing-wave cavity is reciprocating, we denote $E_{it}$ and $E_{ir}$ as the probe cavity-mode field copropagating and counterpropagating with the coupling field $E_c$, respectively. Figure 3 plots the real and imaginary parts of $\chi$ for the two cases in Fig. 2 in free space. The gray dotted lines in Fig. 3 are that of the two-level case ($Ω_c = 0$). For the atom moving with a velocity $v$ that feels the frequency detuning of $E_{it}$ and $E_c$, as $Δ_c - kv$ and $Δ_c - kv$, and the frequency difference between them is Doppler free, that is the two-photon resonance could be satisfied, then a sharp normal dispersion and a reduced absorption can be shown close to the center of atomic resonance, noted by the red dashed lines, that is the EIT effect. Conversely, however, the moving atom feels the frequency detuning of $E_{ir}$ and $E_c$ as $Δ_c + kv$ and $Δ_c - kv$, respectively, a Doppler shift with $2kv$ is created that is not satisfied with two-photon resonance. Here, the dispersion for the intercavity medium to $E_{ir}$ is anomalous, and the strong absorption replaces the transparency for $E_{ir}$ close to the center of the single-photon resonance, noted by the blue solid lines, which is the key factor that suppress the dark-state polariton in SC. In Fig. 3(b), note that the probe cavity mode neither co- or counterpropagates with the coupling field, and a Gaussian absorption profile is always present based on the Doppler background. The width and depth of the absorption profiles depend on the intensity of the coupling, and their absorptive centers are controlled by the coupling detuning.

Due to the difference in the susceptibility feature for the intracavity medium and the reciprocating cavity-mode fields, the intensity transmission function of the coupled atom-cavity system is given by

$$I = \frac{(1-r)^2 \exp(-α_2l)}{[1 - r \exp(-αl)]^2 + 4r \exp(-αl)\sin^2(\psi/2)}, \quad (4)$$

where $r$ is the reflectivity of the mirrors, $α = (α_1 + α_2)/2$, $α_1$ and $α_2$ are the absorption coefficients of intracavity atoms to $E_{it}$ and $E_{ir}$, respectively; $ψ$ is the cavity-mode round-trip phase shift and is described as

$$ψ = \frac{2\pi(Δ_c - δ_0)}{Δ_{FSR}} + \frac{2\pi l}{λ} (n_1 + n_2 - 2), \quad (5)$$

where $Δ_{FSR}$ is the free spectral region (FSR) of the near confocal cavity and $l$ is the length of the Cs cell in the cavity, $n_{1(2)}$ is the refractive index, and $n_{1(2)} - 1$ is proportional to the dispersion of intracavity atoms to the $E_{it(r)}$.

If $α_2 = α_1$ and $n_2 = n_1$, then Eq. (4) describes the intensity transmission function of the RC. Using a lower atomic velocity distribution, Fig. 4 theoretically compares the normalized cavity transmission spectrum for SC and RC. Under the condition of atomic resonance ($Δ_c = δ_0 = 0$), for RC, a narrow central dark-state peak (intracavity EIT) can be created obviously [see the dashed line (I) in Fig. 4]. For SC, however, the dark-state peak is severely suppressed due to the strong absorption of the intracavity atoms to $E_{ir}$ [see the solid line (II) in Fig. 4]; this coincides qualitatively with the experimental result in Fig. 2. Certainly, if the theoretical analysis is in agreement strictly with the experiment, more concrete parameters and factors should be considered on the modeling of this system.

For SC, when the coupling frequency deviates from the atomic resonance ($Δ_c ≠ 0$), it will not happen at the same position for the EIT effect caused by the intracavity medium copropagating $E_{it}$ and the strong absorption of the counterpropagating $E_{ir}$ due to the Doppler shift of the hot atoms (where the frequency difference is $2Δ_c$), so the intracavity dark-state peak cannot be seen until the absorption effect becomes weakened.

It is found experimentally that the dark-state polariton peak can be generated obviously just when the coupling-field frequency deviates from the atomic resonance beyond $~±100$ MHz within the interval of two normal sides under our experimental conditions. Since the absorption of the reflected cavity mode is still strong when the single-photon detuning $Δ_c$ is smaller than $100$ MHz, the...
intracavity atoms are opaque to the probe mode in this range, which is dependent on the temperature of the intracavity atoms. Figure 5 presents the cavity transmission spectra versus the probe detuning at the different coupling and cavity-mode detunings. When the coupling frequency $\Delta_c$ is negatively detuned larger than 100 MHz at $\delta_q = 0$, a narrow intracavity dark-state polariton peak is created at the position of $\Delta_c = \Delta_q$, and its intensity rising with $\Delta_c$ is increased, see Figs. 5(a1) and 5(b1). The linewidth of the dark-state polariton peak is measured to be ~4 MHz, which is limited primarily by the composite cavity fineness in our experiment, and is smaller than that of the far-off resonance cavity mode (~24 MHz) and the EIT signal in free space (~13 MHz). Furthermore, when $\Delta_c$ is left detuned to superpose with the left normal mode, the intracavity EIT disappears and the left normal mode is splitting two sides [see Fig. 5(c1)]. The dashed lines are the relevant theoretical simulations, which are in reasonably agreement with the experimental results.

When the coupling detuning $\Delta_c$ is locked and shifting the cavity-mode frequency $\delta_q$, the anticrossing-like behavior is shown for the position of the two normal side peaks, and has little influence on that of the dark-state polariton peak[23], but its transmission intensity is sensitive. For the specific $\Delta_c$, the value of $\delta_q$ to achieve the max efficiency of the intracavity EIT peak is different. In Fig. 5(a2), when $\Delta_c = -116$ MHz, the max intracavity EIT is created at $\delta_q = -8$ MHz, just the increase is not obvious compared to the resonant center because of absolute intensity is small. For $\Delta_c = -185$ MHz, however, the max efficiency raises 1.5 times when $\delta_q$ is shifted to ~70 MHz, meanwhile its linewidth remains unchanged, and its position is shifted left about 2 MHz at the large shifting range of $\delta_q$ due to the pulling effect, as shown in Fig. 5(b2).

Surprisingly, the intracavity EIT appears with efficient intensity and narrower linewidth when $\delta_q$ reached -120 MHz [see the solid line in Fig. 5(c2)]. That is because the frequency pulling effect of the cavity-mode detuning to the left normal mode is nearly linear with $\delta_q$ increased, whose frequency position is separated obviously from that of the intracavity EIT. Obviously, the shift value of $\delta_q$ to the maxium of the intracavity EIT is not in line with $\Delta_c$ detuned, which is mainly because the influence of $\delta_q$ on the intensity is asynchronous compared with $\Delta_c$, as in Eq. (4). Note that it has some differences between the dashed line and the solid line in Fig. 5, because the bistability is generated for the normal mode in the experiment[15], but not considered in the theory.

In the experiment, on the basis of a guarantee to create dark-state polariton peaks, a profile with six peaks can be shown within one FSR by properly adjusting the experimental parameters. Figure 6 presents the experimentally measured cavity transmission spectrum at $\Delta_c = -176$ MHz. Comparing with Fig. 5(b2), one can see the similar change for the dark-state polariton peak with $\delta_q$ negatively detuned. In Fig. 6(d), at the resonant case ($\Delta_p = 0$), the central little peak is the residual resonant transmission, the left and right normal splitting modes split into two peaks, respectively. When the cavity mode is negatively shifted, for the left-hand normal mode peak, one of the split peaks moves outward and another moves inward [see Figs. 6(b)–6(d)]. When $\delta_q = -76$ MHz, the inward-moving peak merges with the central peak [see Fig. 6(a)]. For the right-hand normal mode peak, however, the two split peaks close gradually until merging into one peak [see Fig. 6(b)]. When the cavity detuning positively detuned, one can see the contrary result. Note that the influence of the dark-state polariton when $\delta_q = 48$ MHz shows a rather dispersive-

![Fig. 5. The cavity transmission spectra versus the probe detuning at different coupling and cavity-mode detunings.](image-url)

![Fig. 6. Experimentally measured multipk cavity transmission spectrum with $\Delta_c = -176$ MHz for different cavity-mode detunings.](image-url)
like peak, not one merging peak as is shown for the left normal mode [see Fig. 6(f)].

In this Letter, we present our experimental investigations of the cavity transmission spectrum with hot three-level atoms coupled to an optical SC. Different from the result of the RC, the dark-state polariton peak is strongly inhibited at the atomic resonance due to the strong absorption effect of the counterpropagating cavity mode. When the coupling frequency is off resonance, a narrow linewidth dark-state polariton peak is created because the Doppler shift reduces the absorption of the counterpropagating field. Furthermoe, the optimum position of the cavity-mode frequency for the efficient dark-state polariton peak is theoretically and experimentally demonstrated. By changing parameters, a multipeak cavity transmission is observed experimentally. This study helps us to better understand the dynamic process of the coherent atoms coupled to the optical cavity system, which has potential applications in broadband, multichannel quantum information processing, and all-optical quantum devices.

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References